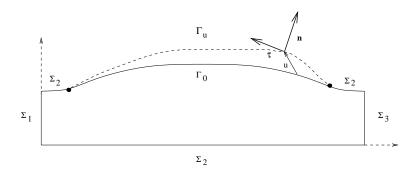
# Numerical control of normal velocity by normal stress for fluid-structure interaction problem

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#### Geometrical model and notations



# Notations from thin shell theory

$$L>0$$
,  $\phi^1,\phi^2:[0,L]\to\mathbb{R}$ 

$$\Gamma_0 = \{(x_1, x_2) \in \mathbb{R}^2; \ x_1 = \phi^1(\xi), \ x_2 = \phi^2(\xi), \ \xi \in [0, L]\}$$

The covariant basis  $(a_1, a_3)$  is given by

$$\mathbf{a_1} = \left( \left( \phi^1 \right)', \left( \phi^2 \right)' \right)^T, \quad \mathbf{a_3} = \frac{1}{\sqrt{a}} \left( \left( \phi^2 \right)', - \left( \phi^1 \right)' \right)^T$$

and the associated contravariant basis  $\left(a^{1},a^{3}\right)$  is

$$\mathbf{a}^1 = \frac{1}{\sqrt{a}}\mathbf{a}_1, \quad \mathbf{a}^3 = \mathbf{a}_3$$

where 
$$a = ((\phi^1)')^2 + ((\phi^2)')^2$$
.

# Weak formulation of the arch problem

$$\mathbf{U} = \left\{ \boldsymbol{\psi} = \left(\psi_1, \psi_3\right) \in H_0^1\left(\left]0, L\right[\right) \times H_0^2\left(\left]0, L\right[\right) \right\}$$

$$\mathbf{a}_{\mathcal{S}}\left(\mathbf{u},\boldsymbol{\psi}\right)=\int_{0}^{L}\frac{\mathit{Ee}}{1-\nu^{2}}\left(\gamma_{1}^{1}\left(\mathbf{u}\right)\gamma_{1}^{1}\left(\boldsymbol{\psi}\right)+\frac{\mathit{e}^{2}}{12}\rho_{1}^{1}\left(\mathbf{u}\right)\rho_{1}^{1}\left(\boldsymbol{\psi}\right)\right)\sqrt{\mathit{a}}\;\mathit{d}\xi$$

Find the displacement  $\mathbf{u}=(u_1,u_3)\in U$  of the arch such that for all  $\psi\in U$ , we have

$$a_{S}(\mathbf{u}, \psi) = \int_{0}^{L} (\lambda^{1} \psi_{1} + \lambda^{3} \psi_{3}) \sqrt{a} d\xi + \int_{0}^{L} (f^{S,1} \psi_{1} + f^{S,3} \psi_{3}) \sqrt{a} d\xi$$

$$\Gamma_u = \left\{ (x_1, x_2)^T = \phi(\xi) + u_1(\xi) \mathbf{a^1}(\xi) + u_3(\xi) \mathbf{a^3}(\xi), \xi \in [0, L] \right\}$$



### Strong form of the fluid equations

Find the velocity  $\mathbf{v}:\Omega_u^F\to\mathbb{R}^2$  and the pressure  $p:\Omega_u^F\to\mathbb{R}$  such that:

$$\begin{array}{rcl} -\mu \Delta \mathbf{v} + \nabla p & = & \mathbf{f}^F, & \text{in } \Omega^F_u \\ \nabla \cdot \mathbf{v} & = & 0, & \text{in } \Omega^F_u \\ \mathbf{v} \times \mathbf{n} & = & 0, & \text{on } \Sigma_1 \\ \mathbf{n} \cdot \left(\sigma^F \mathbf{n}\right) & = & -P_{in}, & \text{on } \Sigma_1 \\ \mathbf{v} & = & 0, & \text{on } \Sigma_2 \\ \mathbf{v} \times \mathbf{n} & = & 0, & \text{on } \Sigma_3 \\ \mathbf{n} \cdot \left(\sigma^F \mathbf{n}\right) & = & -P_{out}, & \text{on } \Sigma_3 \\ \mathbf{v} \times \mathbf{n} & = & 0, & \text{on } \Gamma_u \\ \mathbf{n} \cdot \left(\sigma^F \mathbf{n}\right) & = & -\lambda^3 \circ \left(\phi + u_1 \mathbf{a}^1 + u_3 \mathbf{a}^3\right)^{-1}, & \text{on } \Gamma_u \end{array}$$

# Week formulation of the fluid problem. Notations

$$\mathbf{W} = \left\{ \mathbf{w} \in \left( H^{1} \left( \Omega_{u}^{F} \right) \right)^{2}; \ \mathbf{w} = 0 \text{ on } \Sigma_{2}, \ w_{2} = 0 \text{ on } \Sigma_{1} \cup \Sigma_{3} \right\},$$

$$Q = L^{2} \left( \Omega_{u}^{F} \right), \quad Z = H^{1/2} \left( \Gamma_{u} \right).$$

$$a_{F} \left( \mathbf{v}, \mathbf{w} \right) = 2\mu \sum_{i,j=1}^{2} \int_{\Omega_{u}^{F}} \epsilon_{ij}(\mathbf{v}) \epsilon_{ij}(\mathbf{w}) d\mathbf{x}$$

$$b_{F} \left( \mathbf{w}, q \right) = -\int_{\Omega_{u}^{F}} \left( \nabla \cdot \mathbf{w} \right) q d\mathbf{x}$$

$$c_{F} \left( \mathbf{w}, \zeta \right) = -\int_{\Gamma_{u}} \left( \mathbf{w} \cdot \boldsymbol{\tau} \right) \zeta d\gamma$$

$$\ell_{F} \left( \mathbf{w} \right) = \int_{\Omega_{u}^{F}} \mathbf{f}^{F} \cdot \mathbf{w} d\mathbf{x} - \int_{\Gamma_{u}} \left( \mathbf{w} \cdot \mathbf{n} \right) \lambda^{3} \circ \left( \phi + u_{1} \mathbf{a}^{1} + u_{3} \mathbf{a}^{3} \right)^{-1} d\gamma$$

$$-\int_{\Sigma_{1}} \left( \mathbf{w} \cdot \mathbf{n} \right) P_{in} d\gamma - \int_{\Sigma_{3}} \left( \mathbf{w} \cdot \mathbf{n} \right) P_{out} d\gamma$$

### Week formulation of the fluid problem

Find  $\mathbf{v} \in \mathbf{W}$ ,  $p \in Q$  and  $\eta \in Z$  such that

$$a_{F}(\mathbf{v}, \mathbf{w}) + b_{F}(\mathbf{w}, p) + c_{F}(\mathbf{w}, \eta) = \ell_{F}(\mathbf{w}), \ \forall \mathbf{w} \in \mathbf{W}$$
  
 $b_{F}(\mathbf{v}, q) = 0, \ \forall q \in Q$   
 $c_{F}(\mathbf{v}, \zeta) = 0, \ \forall \zeta \in Z$ 

The physical meaning of  $\eta$  is that of the tangential stresses on the interface, i.e.  $\eta = \boldsymbol{\tau} \cdot \boldsymbol{\sigma}^F \mathbf{n}$ .

# Continuity of the stresses and of the velocity on the interface

The surface forces acting to the arch in a point  $\phi(\xi)$  of  $\Gamma_0$  are given by

$$\lambda^1 \mathbf{a_1} + \lambda^3 \mathbf{a_3},$$

while the surface forces acting to the fluid in the point  $\phi(\xi) + u_1(\xi)\mathbf{a}^1(\xi) + u_3(\xi)\mathbf{a}^3(\xi)$  of  $\Gamma_u$  are given by

$$\sigma^{F}\mathbf{n} = \left(\mathbf{\tau}\cdot\left(\sigma^{F}\mathbf{n}\right)\right)\mathbf{\tau} + \left(\mathbf{n}\cdot\left(\sigma^{F}\mathbf{n}\right)\right)\mathbf{n}.$$

$$\frac{1}{\sqrt{a}}\mathbf{a}_{1} = \mathbf{a}^{1} \approx \mathbf{\tau}, \quad \mathbf{a}_{3} = \mathbf{a}^{3} \approx \mathbf{n}$$

$$-\lambda^{1}\sqrt{\mathbf{a}} \approx \boldsymbol{\tau} \cdot \left(\sigma^{F}\mathbf{n}\right) = \eta, \quad -\lambda^{3} \approx \mathbf{n} \cdot \left(\sigma^{F}\mathbf{n}\right)$$
$$\mathbf{v} \cdot \mathbf{n} = 0, \quad \mathbf{v} \times \mathbf{n} = 0, \text{ on } \Gamma_{u}$$

### Coupling fluid and arch equations

For a given  $\lambda$ , we obtain the displacement  $\mathbf{u}$  by solving the arch problem.

From the fluid problem, we get the velocity  $\mathbf{v}$ , the pressure p of the fluid and the tangential stresses on the interface  $\eta$ .

The fluid-arch interaction problem is to find  $\lambda$  such that

$$\mathbf{v} \cdot \mathbf{n} = 0$$
, on  $\Gamma_u$   
 $-\lambda^1 \sqrt{a} = \eta \circ \left(\phi + u_1 \mathbf{a}^1 + u_3 \mathbf{a}^3\right)$ , on  $[0, L]$ .

### Optimal control problem

We assume in the following that the displacement of the arch depends only on  $\lambda^3$ .

$$\inf \frac{1}{2} \int_{\Gamma_u} (\mathbf{v} \cdot \mathbf{n})^2 \ d\gamma$$

subject to

$$\lambda^3 \in L^2\left(]0,L[
ight)$$
  $u$  solution of the arch problem  $\mathbf{v},\ p,\ \eta$  solution of the fluid problem

#### Geometry

Let L = 6 cm.

$$\Gamma_0 = \left\{ (x_1, x_2) \in \mathbb{R}^2; \ x_1 = \xi, \ x_2 = -5 + \sqrt{45 - (\xi - 3)^2}, \ \xi \in [0, L] \right\}.$$

#### Physical parameters

The inflow  $\Sigma_1$  and outflow  $\Sigma_3$  sections are segments of length 0.8. The thickness of wall a=0.1 cm, the Young's modulus

The thickness of wall e = 0.1 cm, the Young's modulus  $E = 0.75 \cdot 10^6$  g the Poisson's ratio v = 0.40 and the

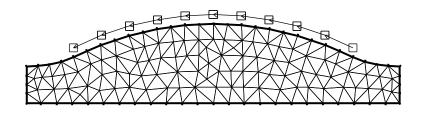
 $E=0.75\cdot 10^6 \frac{g}{cm\cdot s^2}$ , the Poisson's ratio  $\nu=0.49$  and the mass density  $\rho^S=1.1 \frac{g}{cm^3}$ .

The viscosity of the fluid  $\mu = 0.035 \frac{g}{cm \cdot s}$ , the volume force in fluid is  $\mathbf{f}^F = (0,0)^T$  and the outflow pressure  $P_{out} = 0$ .

### Finite Element approximation

The normal displacement of the arch is approached by the finite element  $\mathbb{P}_3$  Hermite, while the tangential displacement is approached by  $\mathbb{P}_1$  on segments.

For the approximation of the fluid velocity and pressure we have employed the triangular finite elements  $\mathbb{P}_1$ +bubble and  $\mathbb{P}_1$  respectively. The tangential stresses on the interface is approached by the finite element  $\mathbb{P}_1$  on segments.



#### Construction of fluid meshes

All the fluid meshes are obtained by moving a fixed mesh with a displacement which is the solution of a Laplace problem with Dirichlet boundary conditions.

On the fixed boundaries  $\Sigma_1$ ,  $\Sigma_2$ ,  $\Sigma_3$  the mesh displacement vanishes, while on the interface  $\Gamma_u$ , it is equal to the arch displacement.

### Solving the optimal control problem

In order to solve numerically the optimal control problem, we have used the Broyden, Fletcher, Goldforb, Shano (BFGS) method where the gradient of the cost function was approached by the first order Finite Difference Method with the grid spacing 0.01. We can use the stopping criteria  $|J| < \epsilon_1$  or  $\|\nabla J\|_{\infty} < \epsilon_2$ .

# Fluid mesh with 20 segments on the interface

Iteration	J	Iteration	J
0	2547.091786	8	1.023813629
1	1211.782266	9	0.541353106
2	725.6283591	10	0.196142907
3	3 510.1113951		0.074620024
4	55.08303685	12	0.061075290
5	5 14.69322181		0.059283206
6	6 2.761055109		0.059257208
7	1.355973240		

$$\|\nabla J\|_{\infty} = 1.9 \cdot 10^2$$

# Fluid mesh with 50 segments on the interface

Iteration	J	Iteration	J
0	2961.338390	7	8.851579284
1	2160.756802	8	2.369258721
2	1392.917188	9	1.271752295
3	879.6830449	10	0.382701032
4	206.9249346	11	0.179680309
5	5 28.11455274		0.124469558
6	18.81352198	13	0.123437472

$$\|\nabla J\|_{\infty} = 1.0 \cdot 10^2$$

# Fluid mesh with 80 segments on the interface

Iteration	J	Iteration	J
0	3619.097215	6	29.86861359
1	1503.175625	7	13.36326171
2	258.7993032	8	10.33965940
3	91.91149097	9	0.753294422
4	61.75894998	10	0.095747979
5	42.29953564	11	0.092956444

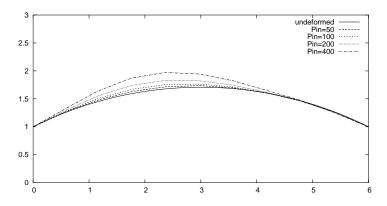
$$\|\nabla J\|_{\infty} = 3.3 \cdot 10^2$$

# Results obtained by BFGS method for different inflow pressure

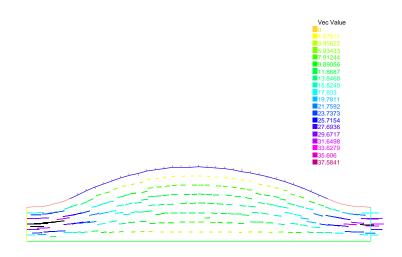
The arch mesh has 10 segments. All the fluid meshes have 516 triangles, 304 vertices and 30 segments on the elastic boundary  $\Gamma_u$ . We have solved numerically the fluid-structure problem for following inflow pressure  $P_{in}=50,\ 100,\ 200,\ 400\ dyn/cm^2$ . The stopping criteria was: |J|<0.5 or  $\|\nabla J\|_{\infty}<0.1$ .

Pin	initial <i>J</i>	final J	no. BFGS iterations
50	2928.63	0.23	9
100	356.86	0.26	13
200	1378.34	0.56	17
400	5257.25	2.08	20

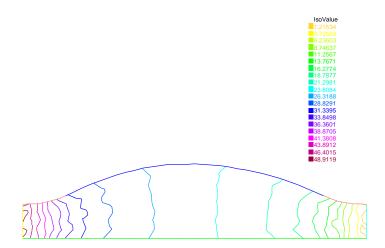
#### Arch deformations for different inflow pressures



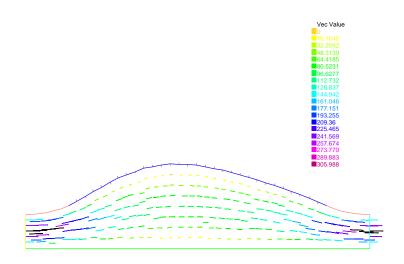
# Fluid velocities scaled by a factor 0.008 in the case $P_{in} = 50$



# Fluid pressure in the case $P_{in} = 50$



# Fluid velocities scaled by a factor 0.001 in the case $P_{in} = 400$



# Fluid pressure in the case $P_{in} = 400$

