

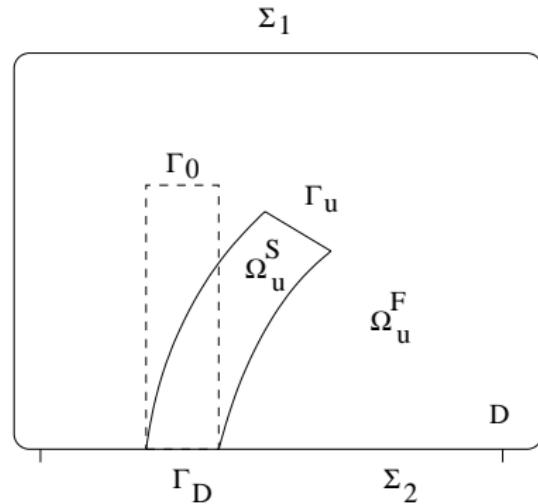
# Approximation of a fluid-structure interaction problem using fictitious domain approach with penalization

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# Geometrical configuration



$$\mathbf{x} = \varphi(\mathbf{X}) = \mathbf{X} + \mathbf{u}(\mathbf{X})$$

The deformed domain  $\Omega_u^S = \varphi(\Omega_0^S)$ .

## Strong formulation

Find  $\mathbf{u} : \overline{\Omega}_0^S \rightarrow \mathbb{R}^2$ ,  $\mathbf{v} : \overline{\Omega}_u^F \rightarrow \mathbb{R}^2$  and  $p : \overline{\Omega}_u^F \rightarrow \mathbb{R}$  such that

$$-\nabla_{\mathbf{x}} \cdot \sigma^S(\mathbf{u}) = \mathbf{f}^S, \quad \text{in } \Omega_0^S \quad (1)$$

$$\mathbf{u} = 0, \quad \text{on } \Gamma_D \quad (2)$$

$$-\nabla \cdot \sigma^F(\mathbf{v}, p) = \mathbf{f}^F, \quad \text{in } \Omega_u^F \quad (3)$$

$$\nabla \cdot \mathbf{v} = 0, \quad \text{in } \Omega_u^F \quad (4)$$

$$\mathbf{v} = \mathbf{g}, \quad \text{on } \Sigma_1 \quad (5)$$

$$\mathbf{v} = 0, \quad \text{on } \Sigma_2 \setminus \Gamma_D \quad (6)$$

$$\mathbf{v} = 0, \quad \text{on } \Gamma_u \quad (7)$$

$$\omega \left( \sigma^F(\mathbf{v}, p) \mathbf{n}^F \right) \circ \varphi = -\sigma^S(\mathbf{u}) \mathbf{n}^S, \quad \text{on } \Gamma_0 \quad (8)$$

## Fictitious domain approach using penalization

$$\begin{aligned}-\nabla \cdot \sigma^F(\mathbf{v}, p) + \frac{1}{\epsilon} \mathcal{P}(\mathbf{v}) &= \mathbf{f}^F, \quad \text{in } \Omega_u^S \\ \nabla \cdot \mathbf{v} &= 0, \quad \text{in } \Omega_u^S\end{aligned}$$

where  $\epsilon > 0$  is a penalization parameter,

$$\mathcal{P}(\mathbf{v}) = \left( |v_1|^{\alpha-1} \operatorname{sgn}(v_1), |v_2|^{\alpha-1} \operatorname{sgn}(v_2) \right)$$

where  $\mathbf{v} = (v_1, v_2)$  and  $1 < \alpha < 2$  is a real number.

$$\chi_u^S(\mathbf{x}) = \begin{cases} 1, & \mathbf{x} \in \overline{\Omega}_u^S \\ 0, & \mathbf{x} \in D \setminus \overline{\Omega}_u^S \end{cases} \quad \text{and} \quad \chi_u^F = 1 - \chi_u^S.$$

$$\begin{aligned}-\nabla \cdot \sigma^F(\mathbf{v}, p) + \frac{1}{\epsilon} \chi_u^S \mathcal{P}(\mathbf{v}) &= \mathbf{f}^F, \quad \text{in } D \\ \nabla \cdot \mathbf{v} &= 0, \quad \text{in } D.\end{aligned}$$

# Weak formulation in the structure domain (I)

## Structure equations in the initial domain

$$\int_{\Omega_0^S} \sigma^S(\mathbf{u}) : \nabla_{\mathbf{X}} \mathbf{w}^S d\mathbf{X} = \int_{\Omega_0^S} \mathbf{f}^S \cdot \mathbf{w}^S d\mathbf{X} + \int_{\Gamma_0} \sigma^S(\mathbf{u}) \mathbf{n}^S \cdot \mathbf{w}^S dS.$$

## Fluid equations in the deformed structure domain

We define  $\tilde{\mathbf{w}}^S : \Omega_u^S \rightarrow \mathbb{R}^2$ ,  $\tilde{\mathbf{w}}^S = \mathbf{w}^S \circ \varphi^{-1}$

$$\begin{aligned} & \int_{\Omega_u^S} \sigma^F(\mathbf{v}_\epsilon, p_\epsilon) : \nabla \tilde{\mathbf{w}}^S d\mathbf{x} + \frac{1}{\epsilon} \int_{\Omega_u^S} \mathcal{P}(\mathbf{v}_\epsilon) \cdot \tilde{\mathbf{w}}^S d\mathbf{x} \\ &= \int_{\Omega_u^S} \mathbf{f}^F \cdot \tilde{\mathbf{w}}^S d\mathbf{x} - \int_{\Gamma_u} \sigma^F(\mathbf{v}_\epsilon, p_\epsilon) \mathbf{n}^F \cdot \tilde{\mathbf{w}}^S ds \end{aligned}$$

## Weak formulation in the structure domain (II)

### Fluid equations in the undeformed structure domain

$$\begin{aligned} & \int_{\Omega_0^S} J \left( \sigma^F (\mathbf{v}_\epsilon, p_\epsilon) \circ \varphi \right) \mathbf{F}^{-T} : \nabla_{\mathbf{X}} \mathbf{w}^S d\mathbf{X} + \frac{1}{\epsilon} \int_{\Omega_0^S} J \mathcal{P} (\mathbf{v}_\epsilon \circ \varphi) \cdot \mathbf{w}^S d\mathbf{X} \\ &= \int_{\Omega_0^S} J \left( \mathbf{f}^F \circ \varphi \right) \cdot \mathbf{w}^S d\mathbf{X} - \int_{\Gamma_0} \omega \left( \sigma^F (\mathbf{v}_\epsilon, p_\epsilon) \mathbf{n}^F \circ \varphi \right) \cdot \mathbf{w}^S dS. \end{aligned}$$

Subtracting from the structure equations, we get

$$\begin{aligned} & \int_{\Omega_0^S} \sigma^S (\mathbf{u}) : \nabla_{\mathbf{X}} \mathbf{w}^S d\mathbf{X} - \int_{\Omega_0^S} \mathbf{f}^S \cdot \mathbf{w}^S d\mathbf{X} \\ &= \int_{\Omega_0^S} J \left( \sigma^F (\mathbf{v}_\epsilon, p_\epsilon) \circ \varphi \right) \mathbf{F}^{-T} : \nabla_{\mathbf{X}} \mathbf{w}^S d\mathbf{X} \\ &+ \frac{1}{\epsilon} \int_{\Omega_0^S} J \mathcal{P} (\mathbf{v}_\epsilon \circ \varphi) \cdot \mathbf{w}^S d\mathbf{X} - \int_{\Omega_0^S} J \left( \mathbf{f}^F \circ \varphi \right) \cdot \mathbf{w}^S d\mathbf{X} \end{aligned}$$

# Constitutive relations

$$\epsilon(\mathbf{w}) = \frac{1}{2} \left( \nabla \mathbf{w} + (\nabla \mathbf{w})^T \right).$$

Linear elasticity

$$a_S(\mathbf{u}, \mathbf{w}^S) = \int_{\Omega_0^S} \left( \lambda^S (\nabla \cdot \mathbf{u}) (\nabla \cdot \mathbf{w}^S) + 2\mu^S \epsilon(\mathbf{u}) : \epsilon(\mathbf{w}^S) \right) d\mathbf{x}.$$

Stokes equations

$$\begin{aligned} a_F(\mathbf{v}, \mathbf{w}) &= \int_D 2\mu^F \epsilon(\mathbf{v}) : \epsilon(\mathbf{w}) d\mathbf{x} \\ b_F(\mathbf{w}, p) &= - \int_D (\nabla \cdot \mathbf{w}) p d\mathbf{x}. \end{aligned}$$

## Parametrization of the characteristic function

Let  $j \in W^{1,\infty}(D)$  be a parametrization of  $\Omega_0^S \subset D$ , i.e. :

$$\begin{aligned} j(\mathbf{x}) &> 0, \quad \mathbf{x} \in \Omega_0^S, \\ j(\mathbf{x}) &< 0, \quad \mathbf{x} \in D \setminus \overline{\Omega}_0^S, \\ j(\mathbf{x}) &= 0, \quad \mathbf{x} \in \partial\Omega_0^S. \end{aligned}$$

$\Omega_u^S = \varphi(\Omega_0^S)$ , where  $\varphi(\mathbf{X}) = \mathbf{X} + \mathbf{u}(\mathbf{X})$

$$j_u(\mathbf{y}) = \begin{cases} j(\mathbf{x}), & \mathbf{y} = \varphi(\mathbf{x}) \in \Omega_u^S \\ 0, & \mathbf{y} \in \partial\Omega_u^S \\ -\text{dist}(\mathbf{y}, \overline{\Omega}_u^S), & \mathbf{y} \notin \overline{\Omega}_u^S \end{cases}$$

is a parametrization of  $\Omega_u^S$ ,  $j_u \in W^{1,\infty}(D)$ .

## Regularization of the characteristic function

If  $H$  is the Heaviside function  $H : \mathbb{R} \rightarrow \{0, 1\}$ ,

$$H(r) = \begin{cases} 1, & r \geq 0 \\ 0, & r < 0 \end{cases}$$

then  $H(j_u(\cdot))$  is the characteristic function of  $\Omega_u^S$ .

$$\Omega_0^\epsilon \subset \subset \Omega_0^S, \quad \Omega_u^\epsilon = (\mathbf{id} + \mathbf{u})(\Omega_0^\epsilon)$$

There exists  $\mu_\epsilon > 0$  such that  $j(\mathbf{x}) \geq \mu_\epsilon > 0$ , for all  $\mathbf{x} \in \Omega_0^\epsilon$ .  
Then we take  $\tilde{H} = H^{\mu_\epsilon}$ , the Yosida regularization of  $H$

$$H^{\mu_\epsilon}(r) = \begin{cases} 1, & r \geq \mu_\epsilon \\ \frac{r}{\mu_\epsilon}, & 0 \leq r < \mu_\epsilon \\ 0, & r < 0 \end{cases}$$

It follows that  $H^{\mu_\epsilon}(j_u(\mathbf{x})) = 1$  for all  $\mathbf{x} \in \Omega_u^\epsilon$ .

$\tilde{H}(j_u)$  is Lipschitz and  $0 \leq \tilde{H}(j_u(\mathbf{x})) \leq 1$  for all  $\mathbf{x} \in \overline{D}$ .

# Functional spaces

$$W^S = \left\{ \mathbf{w}^S \in \left( H^1(\Omega_0^S) \right)^2 ; \mathbf{w}^S = 0 \text{ on } \Gamma_D \right\},$$

$$W = (H_0^1(D))^2,$$

$$Q = L_0^2(D) = \{q \in L^2(D); \int_D q \, dx = 0\}.$$

$\mathbf{f}^F \in (L^2(D))^2$ ,  $\mathbf{f}^S \in (L^2(\Omega_0^S))^2$ ,  $\mathbf{g} \in (H^{1/2}(\partial D))^2$ , such that  
 $\mathbf{g} = 0$  on  $\Sigma_2$  and  $\int_{\Sigma_1} \mathbf{g} \cdot \mathbf{n}^F ds = 0$ .

For a given  $\mathbf{u} \in (W^{1,\infty}(\Omega_0^S))^2$ , such that  $\|\mathbf{u}\|_{1,\infty,\Omega_0^S} < 1$  and  $\mathbf{u} = 0$  on  $\Gamma_D$ , we define:

- ▶ fluid velocity  $\mathbf{v}_\epsilon \in (H^1(D))^2$ ,  $\mathbf{v}_\epsilon = \mathbf{g}$  on  $\partial D$ ,
- ▶ fluid pressure  $p_\epsilon \in Q$ ,
- ▶ structure displacement  $\mathbf{u}_\epsilon \in W^S$

## Coupled system

$$a_F(\mathbf{v}_\epsilon, \mathbf{w}) + b_F(\mathbf{w}, p_\epsilon) \\ + \frac{1}{\epsilon} \int_D \tilde{H}(j_u) \mathcal{P}(\mathbf{v}_\epsilon) \cdot \mathbf{w} \, d\mathbf{x} = \int_D \mathbf{f}^F \cdot \mathbf{w} \, d\mathbf{x}, \quad \forall \mathbf{w} \in W \quad (9)$$

$$b_F(\mathbf{v}_\epsilon, q) = 0, \quad \forall q \in Q \quad (10)$$

$$a_S(\mathbf{u}_\epsilon, \mathbf{w}^S) = \int_{\Omega_0^S} \mathbf{f}^S \cdot \mathbf{w}^S \, d\mathbf{X} \\ + \int_{\Omega_0^S} J(\sigma^F(\mathbf{v}_\epsilon, p_\epsilon) \circ \varphi) \mathbf{F}^{-T} : \nabla_{\mathbf{X}} \mathbf{w}^S \, d\mathbf{X} \\ + \frac{1}{\epsilon} \int_{\Omega_0^S} J \tilde{H}(j_u \circ \varphi) \mathcal{P}(\mathbf{v}_\epsilon \circ \varphi) \cdot \mathbf{w}^S \, d\mathbf{X} \\ - \int_{\Omega_0^S} J(\mathbf{f}^F \circ \varphi) \cdot \mathbf{w}^S \, d\mathbf{X}, \quad \forall \mathbf{w}^S \in W^S \quad (11)$$

where  $\varphi(\mathbf{X}) = \mathbf{X} + \mathbf{u}(\mathbf{X})$ ,  $\mathbf{F}(\mathbf{X}) = \mathbf{I} + \nabla_{\mathbf{X}} \mathbf{u}(\mathbf{X})$ ,  $J(\mathbf{X}) = \det \mathbf{F}(\mathbf{X})$ .

## Definition of the nonlinear operator

$$\left\{ \mathbf{u} \in \left( W^{1,\infty}(\Omega_0^S) \right)^2 ; \quad \| \mathbf{u} \|_{1,\infty,\Omega_0^S} < 1, \quad \mathbf{u} = 0 \text{ on } \Gamma_D \right\}$$

Fluid problem

$$\mathbf{u} \rightarrow \mathbf{v}_\epsilon \in \widetilde{\mathbf{g}} + \left( H_0^1(D) \right)^2 \text{ and } p_\epsilon \in L_0^2(D)$$

Structure problem

$$\mathbf{u}, \quad \mathbf{v}_\epsilon, \quad p_\epsilon \rightarrow \mathbf{u}_\epsilon \in \left\{ \mathbf{w}^S \in \left( H^1 \left( \Omega_0^S \right) \right)^2 ; \quad \mathbf{w}^S = 0 \text{ on } \Gamma_D \right\}$$

Nonlinear operator

$$T_\epsilon(\mathbf{u}) = \mathbf{u}_\epsilon.$$

# Existence of the penalized fluid-structure problem

$$B_\delta = \{u \in W^{1,\infty}(\Omega_0^S)^2; \|u\|_{1,\infty,\Omega_0^S} \leq \eta_\delta, u = 0 \text{ on } \Gamma_D\}$$

## Theorem

Let  $\epsilon > 0$  fixed. If  $f^F$ ,  $f^S$  and  $g$  are “small” in their own norms, then the nonlinear operator  $T_\epsilon$  has at least one fixed point  $\mathbf{u}_\epsilon$  in  $B_\delta$ .

See: A. Halanay, C.M. Murea, D. Tiba, Existence and approximation for a steady fluid-structure interaction problem using fictitious domain approach with penalization, in progress

Problem: What happens when  $\epsilon \rightarrow 0$ ?

## Partitioned procedures based on the fixed point iterations

The penalized term  $\mathcal{P}(\mathbf{v}) = \left( |v_1|^{\alpha-1} \operatorname{sgn}(v_1), |v_2|^{\alpha-1} \operatorname{sgn}(v_2) \right)$ , where  $1 < \alpha < 2$  is non-linear in  $\mathbf{v}$ . But, if  $\alpha$  is close to 2, we can approach  $\mathcal{P}(\mathbf{v})$  by  $\mathbf{v}$ . We can also replace  $\tilde{H}(j_u)$  by the characteristic function  $\chi_u^S$ .

### Algorithm

**Step 1.** Given the initial displacement of the structure  $\mathbf{u}^0 \in W^S$ , compute the characteristic function  $\chi_{u^0}^S$ , put  $k := 0$

**Step 2.** Find the velocity  $\mathbf{v}_\epsilon \in (H^1(D))^2$ ,  $\mathbf{v}_\epsilon = \mathbf{g}$  on  $\Sigma_1$ ,  $\mathbf{v}_\epsilon = 0$  on  $\Sigma_2$  and the pressure  $p_\epsilon^k \in Q$  by solving the fluid problem

$$a_F(\mathbf{v}_\epsilon^k, \mathbf{w}) + b_F(\mathbf{w}, p_\epsilon^k) + \frac{1}{\epsilon} \int_D \chi_{u_\epsilon^k}^S \mathbf{v}_\epsilon^k \cdot \mathbf{w} \, d\mathbf{x} = \int_D \mathbf{f}^F \cdot \mathbf{w} \, d\mathbf{x}, \quad \forall \mathbf{w} \in W$$
$$b_F(\mathbf{v}_\epsilon^k, q) = 0, \quad \forall q \in Q$$

**Step 3.** Find the new displacement of the structure  $\mathbf{u}_\epsilon^{k+1} \in W^S$  by solving

$$\begin{aligned} a_S(\mathbf{u}_\epsilon^{k+1}, \mathbf{w}^S) &= \int_{\Omega_0^S} (\mathbf{f}^S - \mathbf{f}^F) \cdot \mathbf{w}^S \, d\mathbf{x} + \int_{\Omega_0^S} 2\mu^F \epsilon(\mathbf{v}_\epsilon^k) : \epsilon(\mathbf{w}^S) \, d\mathbf{x} \\ &\quad - \int_{\Omega_0^S} (\nabla \cdot \mathbf{w}^S) p_\epsilon^k \, d\mathbf{x} + \frac{1}{\epsilon} \int_{\Omega_0^S} (\mathbf{v}_\epsilon^k \circ \varphi_\epsilon^k) \cdot \mathbf{w}^S \, d\mathbf{x} \quad \forall \mathbf{w}^S \in W^S \end{aligned}$$

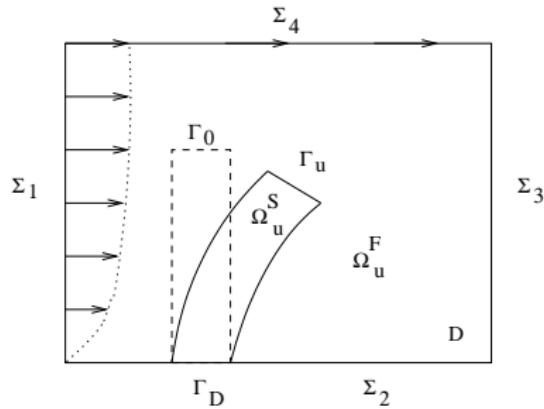
where  $\varphi_\epsilon^k(\mathbf{X}) = \mathbf{X} + \mathbf{u}_\epsilon^k(\mathbf{X})$ .

**Step 4.** Stopping test: if  $\|\mathbf{u}_\epsilon^k - \mathbf{u}_\epsilon^{k+1}\|_{0, \Omega_0^S} \leq tol$ , then **Stop**

**Step 5.** Compute the characteristic function  $\chi_{\mathbf{u}_\epsilon^{k+1}}^S$ , put  $k := k + 1$  and **Go to Step 2.**

Under the assumption of small displacements for the structure, we can approach the Jacobian determinant  $J$  by 1 and the gradient of the deformation  $\mathbf{F}$  by the identity matrix  $\mathbf{I}$ .

# Numerical test 1. Shell in steady-state cross flow



## Physical parameters

Structure: length  $\ell = 3 \text{ m}$ , thickness  $h = 0.125 \text{ m}$ , Young modulus  $E^S = 1.6 \times 10^6 \text{ N/m}^2$ , Poisson's ratio  $\nu^S = 0.49$

Fluid: length  $L = 12 \text{ m}$ , width  $H = 5 \text{ m}$ , mass density  $\rho^F = 1000 \text{ Kg/m}^3$ , dynamic viscosity  $\mu^F = 0.1 \text{ N} \cdot \text{s/m}^2$ . The velocity boundary conditions at the inflow:

$$v_1(x_1, x_2) = V \times 1.5 \frac{(2Hx_2 - x_2^2)}{H^2} \text{ m/s}, \quad V = 5, \quad v_2(x_1, x_2) = 0;$$

at the bottom: no-slip; at the top: slip; at the outflow: the tangential velocity and the normal traction are zero

$$Re = \frac{V\rho^F H}{\mu^F} = 250000$$

## Numerical parameters

Fluid mesh: 82512 triangles and 41682 vertices

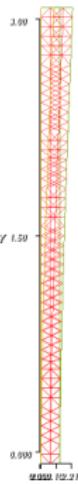
Structure mesh: 188 triangles and 145 vertices

Finite elements: fluid velocity  $\mathbb{P}_1 + \text{bubble}$ , fluid pressure  $\mathbb{P}_1$ , structure displacement  $\mathbb{P}_1$ , characteristic function  $\mathbb{P}_0$

$$\epsilon = 10^{-3}, \quad tol = 10^{-8}$$

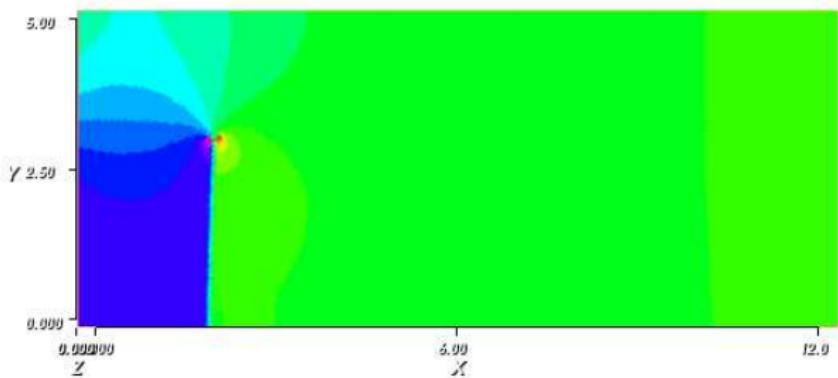
No of iterations of the fixed point algorithm: 7

# The initial and the deformed structure mesh



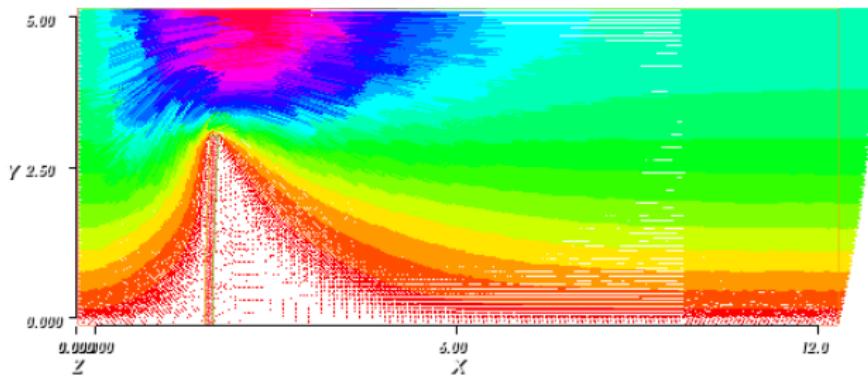
The maximal displacement: 0.083 m

# The fluid pressure

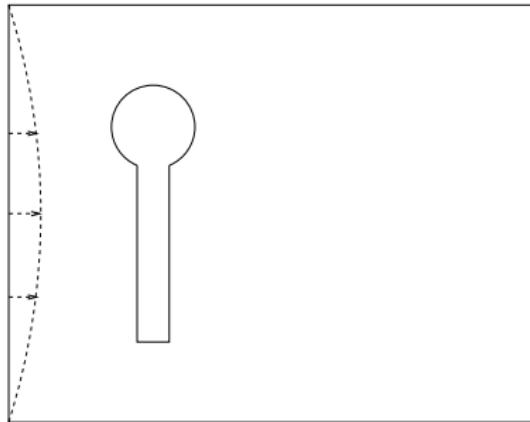


## The fluid velocity around the final position of the structure

$$\|\mathbf{v}_\epsilon\|_{0,\Omega_{u_\epsilon}^S} = \sqrt{\int_D \chi_{u_\epsilon}^S \mathbf{v}_\epsilon \cdot \mathbf{v}_\epsilon d\mathbf{x}} = 0.047270$$



## Numerical test 2. Flexible appendix in a flow



A rectangular flexible appendix is attached to a fixed circle. The circle center is positioned at  $(0.2, 0.2)$  m measured from the left top corner of the channel.

## Physical parameters

Structure: circle radius  $r = 0.5 \text{ m}$ , length  $\ell = 0.35 \text{ m}$ , thickness  $h = 0.02 \text{ m}$ , Young modulus  $E^S = 1.6 \times 10^6 \text{ N/m}^2$ , Poisson's ratio  $\nu^S = 0.49$

Fluid: length  $L = 2.5 \text{ m}$ , width  $H = 0.75 \text{ m}$ , mass density  $\rho^F = 1260 \text{ Kg/m}^3$ , dynamic viscosity  $\mu^F = 1.420 \text{ N} \cdot \text{s/m}^2$ . The velocity boundary conditions at the inflow:

$$v_1(x_1, x_2) = V \times 1.5 \frac{(Hx_2 - x_2^2)}{(H/2)^2} \text{ m/s}, \quad V = 1, \quad v_2(x_1, x_2) = 0;$$

at the bottom and the top: no-slip; at the outflow: traction free

$$Re = \frac{V\rho^F H}{\mu^F} = 665$$

## Numerical parameters

Fluid mesh: 30330 triangles and 15461 vertices

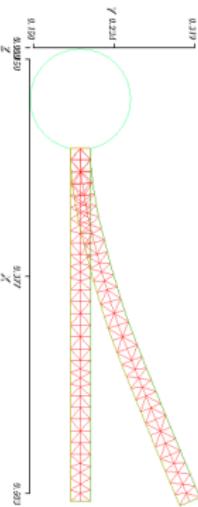
Structure mesh: 128 triangles and 97 vertices

Finite elements: fluid velocity  $\mathbb{P}_1 + bubble$ , fluid pressure  $\mathbb{P}_1$ , structure displacement  $\mathbb{P}_1$ , characteristic function  $\mathbb{P}_0$

$$\epsilon = 10^{-4}, \quad tol = 10^{-8}$$

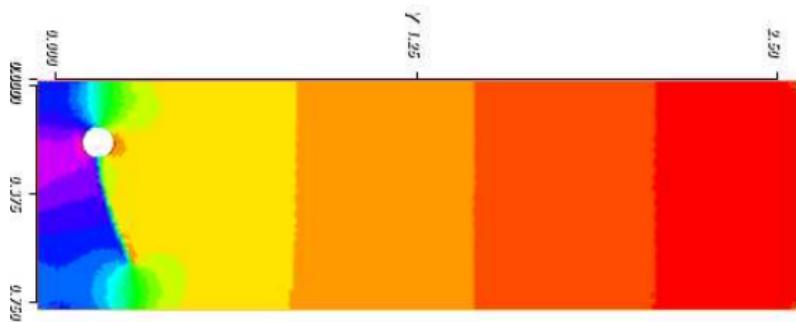
No of iterations of the fixed point algorithm: 8

# The initial and the deformed structure mesh



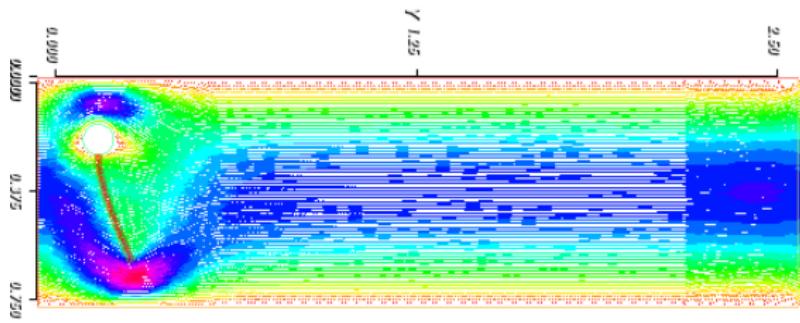
The maximal horizontal displacement: 0.10886 m

# The fluid pressure



# The fluid velocity around the final position of the structure

$$\|\mathbf{v}_\epsilon\|_{0,\Omega_{u_\epsilon}^S} = \sqrt{\int_D \chi_{u_\epsilon}^S \mathbf{v}_\epsilon \cdot \mathbf{v}_\epsilon d\mathbf{x}} = 0.080920$$



Thank you!